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## THE MOLLIER $h-s$ DIAGRAM

### Statement

Analyse the  $h-s$  (Mollier) diagram for a pure substance. In particular:

- Demonstrate that isobars are straight lines in the two-phase region.
- Compute the slope of the saturated vapour curve at low pressure and compare it with data from the Mollier diagram for water.
- Compute the change in slope of isobars at each side of the saturated vapour curve.
- Compute the change in slope of isotherms at each side of the saturated vapour curve.
- How can specific-volume values be computed from isobars?

Se trata del análisis teórico del diagrama  $h-s$  de una sustancia pura. Se pide:

- Demostrar que las isobaras son rectas en la región bifásica.
- Calcular la pendiente de la línea de vapor saturado a bajas presiones y compararla con la del diagrama de Mollier del agua.
- Calcular la variación de pendiente de las isobaras a un lado y otro de la curva de vapor saturado.
- Calcular la variación de pendiente de las isoterms a un lado y otro de la curva de vapor saturado.
- ¿Cómo se calculan volúmenes específicos a partir de las isobaras?

### Solution.

- Demonstrate that isobars are straight lines in the two-phase region.

From  $dh=Tds+vd p$  we deduce that the slope of isobars is:

$$\left. \frac{dh}{ds} \right|_p = T$$

But inside the two-phase region the isotherms coincide with the isobars because they are linked by Clapeyron equation,  $p=p_v(T)$ , given by:

$$\left. \frac{dp}{dT} \right|_{sat} = \frac{s_v - s_L}{v_v - v_L} = \frac{h_v - h_L}{T(v_v - v_L)} = \frac{h_{LV}}{T v_{LV}}$$

thus, the slope of an isobar only depends on its pressure, that is kept constant.

- Compute the slope of the saturated vapour curve at low pressure and compare it with data from the Mollier diagram for water.

From  $dh=Tds+vd p$  we deduce that  $\partial h/\partial s|_{sat}=T+v\partial p/\partial s|_{sat}$ , from  $ds=(c_p/T)dT-\alpha v dp$  we deduce that  $\partial s/\partial p|_{sat}=(c_p/T)\partial T/\partial p|_{sat}-\alpha v$ , and, with  $\partial p/\partial T|_{sat}=h_{LV}/(T v_{LV})$  and the ideal gas approximation we get:

$$\left. \frac{dh}{ds} \right|_{sat} = T + \frac{v}{\left. \frac{ds}{dp} \right|_{sat}} = T + \frac{v}{\frac{c_p}{T} \left. \frac{dT}{dp} \right|_{sat} - \alpha v} = T + \frac{v}{\frac{c_p}{T} \frac{T v_{LV}}{h_{LV}} - \alpha v} = \frac{c_p \frac{T v_{LV}}{h_{LV}} - T \alpha v + v}{\frac{c_p}{T} \frac{T v_{LV}}{h_{LV}} - \alpha v} \stackrel{p \rightarrow 0}{=} \frac{c_p T^2}{c_p T - h_{LV}}$$

With  $c_p=1900 \text{ J/(kg}\cdot\text{K)}$  and  $h_{LV}=2.26\cdot 10^6 \text{ J/kg}$  at  $T=373 \text{ K}$ , we get  $\partial h/\partial s|_{\text{sat}}=-170 \text{ K}$ , whereas by graphical computation from the Mollier diagram for water we get  $\partial h/\partial s|_{\text{sat}}=-140000/1000=-140 \text{ K}$ , not too astray.

- c) Compute the change in slope of isobars at each side of the saturated vapour curve.

It has been deduced that  $\partial h/\partial s|_p=T$ , thus, because isotherms are continuous (not their slopes), the slope of isobars are continuous when traversing the saturated line.

- d) Compute the change in slope of isotherms at each side of the saturated vapour curve.

Similarly:

$$\left. \frac{dh}{ds} \right|_T = T + \frac{v}{\left. \frac{ds}{dp} \right|_T} = T + \frac{v}{-\alpha v} = T - \frac{1}{\alpha}$$

that can be deduced also from the general expressions of  $dh$  and  $ds$  in terms of  $T$  and  $p$ :

$$\left. \frac{dh}{ds} \right|_T = \frac{c_p dT + (1-\alpha T) v dp}{\frac{c_p}{T} dT - \alpha v dp} \Bigg|_T = \frac{(1-\alpha T) v}{-\alpha v} = T - \frac{1}{\alpha}$$

Since  $\alpha$  is discontinuous ( $\alpha=0$  for ideal liquids,  $\alpha=\infty$  for two-phase mixtures,  $\alpha=1/T$  for ideal gases), the slope of the isotherms too. In fact, with the incompressible liquid model the slope becomes infinite, whereas with the ideal gas model the slope is zero.

- e) How can specific-volume values be computed from isobars?

From  $dh=Tds+vdp$  we deduce that  $v=\partial h/\partial p|_s$ .

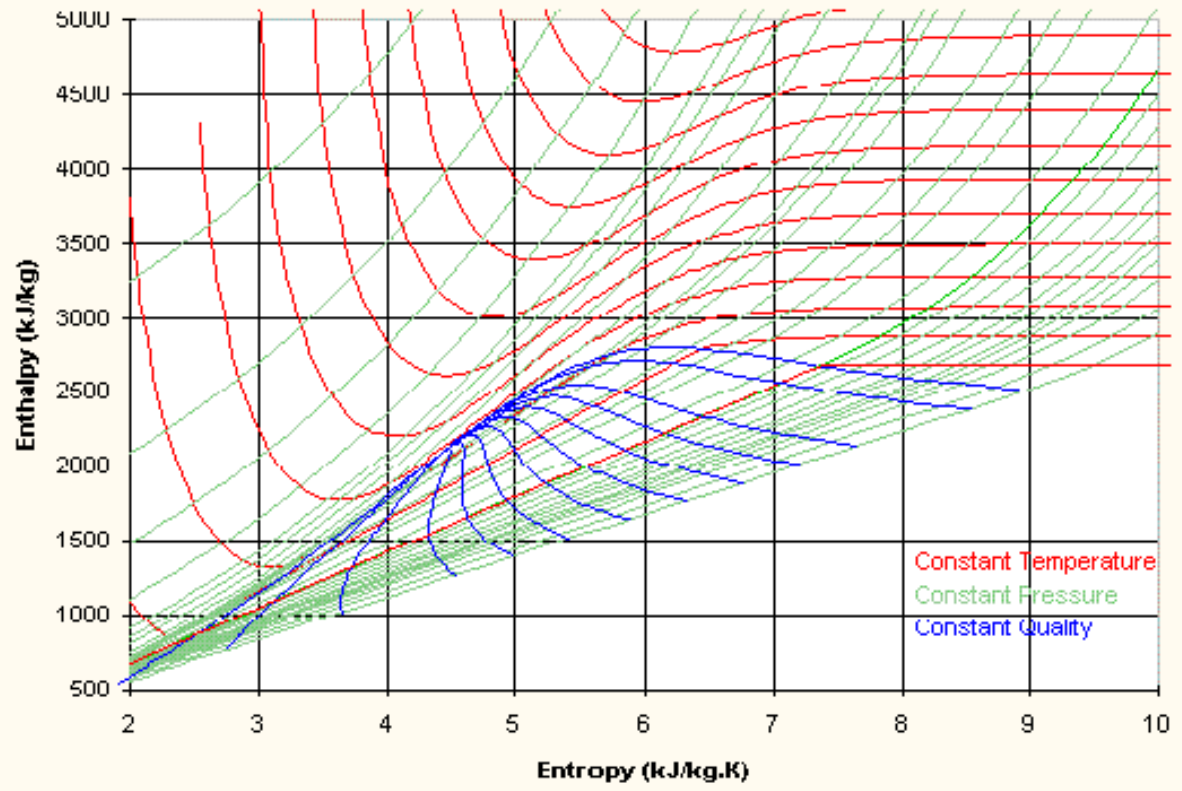
The slope of the isochors is:

$$\left. \frac{dh}{ds} \right|_v = T + \frac{v}{\left. \frac{ds}{dp} \right|_v} = T + \frac{v}{\frac{c_p}{T} \frac{\kappa}{\alpha} - \alpha v} = T + \frac{1}{\frac{\kappa c_v}{\alpha v T}} = T + \frac{\alpha v T}{\kappa c_v} \xrightarrow{\text{PGM}} \left. \frac{dh}{ds} \right|_v = \gamma T$$

Notice that with the ideal gas model, both, isobars and isochors, are exponential functions in the  $h$ - $s$  diagram because their slope grows with  $T$  and  $h$  does, but isochors have a  $\gamma$ -times greater slope.

**Comments.** You may check the results with an actual  $h$ - $s$  diagram as the following one. Notice that the perfect gas model (PGM) only applies to the right of the diagram (i.e. at low pressure).

Steam97 - Mollier Chart



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